

SOME CONDITIONS FOR EXISTENCE AND INTEGRABILITY OF THE FOURIER TRANSFORM

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ABSTRACT. The Fourier transform is naturally defined for integrable functions. Otherwise, it should be stipulated in which sense the Fourier transform is understood. We consider some class of radial and, generally saying, nonintegrable functions. The Fourier transform is calculated as an improper integral and the limit coincides with the Fourier transform in the distributional sense. The inverse Fourier formula is proved as well. Given are some applications of the result obtained.

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1. INTRODUCTION.

1.1. Let \mathbb{R}^n be a real Euclidean n -dimensional space with elements $x = (x_1, \dots, x_n)$. If a function $f(x)$ is integrable, in the Lebesgue sense, over all \mathbb{R}^n there is no problem to define its Fourier transform

$$\hat{f}(x) = \int_{\mathbb{R}^n} f(u) e^{-ix \cdot u} du$$

where $x \cdot u = x_1 u_1 + \dots + x_n u_n$ is the scalar product of $u, x \in \mathbb{R}^n$. But, even in this case, special additional conditions are needed for the inverse formula to be true:

$$f(u) = (2\pi)^{-n} \int_{\mathbb{R}^n} \hat{f}(x) e^{iu \cdot x} dx.$$

For instance, sometimes a summability method should be applied to impart a certain sense to the last integral. And if f is not integrable both formulas need a special investigation.

Thus, the main purpose of our work is to give certain conditions under which both formulas hold. We will formulate them just after introducing notations needed. Then a separate section will be devoted to auxiliary lemmas. The next section will present the proof of the theorem. After that, several sections will be devoted to applications. These applications were announced in [BL1], but the proofs were given only in the author's Ph.D. thesis (see also [L2] where conditions are more restrictive than in this paper). The first application gives a connection between many-dimensional and one-dimensional Fourier transforms. This allows to reduce some multidimensional problems to easier or known one-dimensional ones. Then we give an asymptotic formula for the Fourier transform of a radial function with convexity type conditions.

1.2. We consider radial functions $f(x) = f_0(|x|)$, where $|x| = (x \cdot x)^{\frac{1}{2}}$ is not necessarily integrable over all \mathbb{R}^n . Some notations are needed to describe precisely the class of functions to be studied. For the sake of convenience and completeness, some other notations, which will be used in proofs, are given here as well.

Let $0 < \alpha \leq \frac{n-1}{2}$ and α^* be the greatest integer smaller than α . Denote by

$$W_\alpha(f_0; t) = \frac{1}{\Gamma(\alpha)} \int_t^\infty f_0(r) (r-t)^{\alpha-1} dr$$

the Weyl integral of fractional order where r and t are real numbers. When $0 < \alpha < 1$

$$f_0^{(\alpha)}(t) = \frac{d}{dt} W_{1-\alpha}(f_0; t)$$

is the Weyl derivative of fractional order. When $\alpha = p + \gamma$ with $p = 1, 2, \dots$, and $0 < \gamma < 1$, then

$$f_0^{(\alpha)}(t) = \frac{d^p}{dt^p} f_0^{(\gamma)}(t).$$

Let

$$R_\alpha(f_0; t) = \frac{1}{\Gamma(\alpha)} \int_0^t f_0(r) (t-r)^{\alpha-1} dr$$

be the Riemann-Liouville fractional integral. All these notions may be found for example in [PE2], Ch.12 (see also [Ca], [SKM], [Tr]).

Let $C[a, b]$ and $C^p[a, b]$ be the classes of continuous functions and having p continuous derivatives, respectively.

Let us introduce two Bessel-type functions:

$$\begin{aligned} Q_\alpha(t) &= \int_0^1 (1-s)^{\alpha-1} s^{\frac{n}{2}} J_{\frac{n}{2}-1}(ts) ds \\ &= \Gamma(\alpha) t^{-\frac{n}{2}-\alpha} R_\alpha(s^{\frac{n}{2}} J_{\frac{n}{2}-1}(s); t) \end{aligned}$$

and

$$q_\alpha(t) = \int_0^1 (1-s)^{\alpha-1} s^{\frac{n}{2}-1} J_{\frac{n}{2}}(ts) ds$$

where J_μ is the Bessel function of first type and order μ .

In what follows $F_\alpha(t) = t^{\frac{n-1}{2}} f_0^{(\alpha)}(t)$. When $\alpha = \frac{n-1}{2}$, for brevity, we will write simply Q, q, F .

We define $\varphi \in S$ when $\varphi \in C^\infty$ and $\varphi(x)$ and its derivatives, which we are allowed to multiply by any polynomial of $|x|$, tend to 0 uniformly as $|x| \rightarrow \infty$. Continuous linear forms on S are called tempered distributions.

We shall denote absolute constants, that is independent of substantial parameters, by the letter C . If a dependence on some parameters is essential, they will be indicated as subscripts. Whenever no confusion can result, we use the same letter for different constants in different places.

1.3. Let us consider a class of radial functions satisfying the following conditions:

- (1) $f_0, \dots, f^{(\alpha^*)}$ are locally absolutely continuous on $(0, \infty)$;
- (2) $\lim_{t \rightarrow \infty} t^p f_0^{(p)}(t) = 0$ for $p = 0, 1, \dots, \alpha^*$;
- (3) $\lim_{t \rightarrow \infty} F_\alpha(t) = 0$;
- (4) F_α is a function of bounded variation on $(0, \infty)$.

The total variation of the function F_α will be denoted by V_{F_α} .

Theorem 1. *Let a function f be radial and satisfy conditions (1) - (4). Then for $|x| > 0$*

$$(5) \quad f(x) = \lim_{A \rightarrow \infty} (2\pi)^{-n} \int_{|u| \leq A} \left(1 - \frac{|u|^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} \hat{f}(u) e^{ix \cdot u} du$$

where

$$(6) \quad \hat{f}(u) = \frac{(2\pi)^{\frac{n}{2}} (-1)^{\alpha^*+1}}{\Gamma(\alpha)} |u|^{1-\frac{n}{2}} \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(|u|t) dt$$

is continuous, tends to zero as $|x| \rightarrow \infty$ and coincides with the Fourier transform \tilde{f} , of the function f , in the distributional sense. Both integrals converge uniformly for $|x| \geq r_0 > 0$.

We have to mention that the case $\alpha = \frac{n-1}{2}$ was considered earlier in [BL1], [BL2]. Some ideas of the proof in [BL2] were used here and also in [L2] where such results are established under superfluous conditions at zero.

1.4. Compare Theorem 1 with some earlier results. S. Bochner in [Bo], §44, considers not only radial functions, but more restrictive conditions are claimed for the radial part of a function (spherical average of a function). Analogously, the radially allows less restrictive smoothness conditions than those in V.A. Ilyin and S. A. Alimov theorems for general spectral expansions (see [IA]). In M. L. Goldman's paper [G] radial functions are considered, with "worse" conditions at infinity and the monotonicity of a given function and its derivatives. Very simple formula somehow similar to (6) may be found in [SKM], Ch.5, Lemma 25.1', but the authors did not care much for sharp assumptions.

2. AUXILIARY LEMMAS.

Lemma 1. *The following asymptotic relation holds*

$$q_\alpha(r) = \Gamma(\alpha) r^{-\alpha} J_{\frac{n}{2}+\alpha}(r) + \zeta_{\alpha,n} r^{-\frac{n}{2}} + O\left(r^{-\alpha-\frac{3}{2}}\right)$$

as $r \rightarrow \infty$, and $\zeta_{\alpha,n}$ are some numbers.

Proof. We have

$$\begin{aligned} q_\alpha(r) &= \sum_{j=0}^M \int_0^1 (1-s)^{\alpha-1+j} (1+s)^j s^{\frac{n}{2}+1} J_{\frac{n}{2}}(rs) ds \\ &\quad + \int_0^1 g_0(s) s^{\frac{n}{2}-1} J_{\frac{n}{2}}(rs) ds \end{aligned}$$

where M is such that $g_0(s) = (1-s)^{\alpha+M}(1+s)^{M+1}$ is smooth enough. Evaluate first the last integral. We need the following properties of the Bessel functions (see e.g., [BE2], §7.2.8(50),(51); §7.13.1(3); §7.12(8)):

$$(7) \quad \frac{d}{dt} [t^{\pm\nu} J_\nu(t)] = \pm t^{\pm\nu} J_{\nu \mp 1}(t);$$

$$(8) \quad J_\nu(t) = \sqrt{\frac{2}{\pi t}} \cos\left(t - \frac{\pi\nu}{2} - \frac{\pi}{4}\right) + \sqrt{\frac{2}{\pi}} \frac{1-4\nu^2}{8} t^{-\frac{3}{2}} \sin\left(t - \frac{\pi\nu}{2} - \frac{\pi}{4}\right)$$

$$+ O\left(t^{-\frac{5}{2}}\right) \quad \text{as } t \rightarrow \infty;$$

$$(9) \quad J_\nu = O(t^\nu) \quad \text{for small } t.$$

Let us integrate by parts, using (7) as follows:

$$\begin{aligned} \int_0^1 g_0(s) s^{\frac{n}{2}-1} J_{\frac{n}{2}}(rs) ds &= \int_0^1 [g_0(s) s^{n-2}] [s^{-\frac{n}{2}+1} J_{\frac{n}{2}}(rs)] ds \\ &= -\frac{1}{r} g_0(s) s^{\frac{n}{2}-1} J_{\frac{n}{2}-1}(rs) \Big|_0^1 + \frac{1}{r} \int_0^1 g_1(s) s^{\frac{n}{2}-2} J_{\frac{n}{2}-1}(rs) ds \end{aligned}$$

where g_1 is also smooth enough. For the case $n = 2$, unlike for that for $n > 2$ the factor $s^{\frac{n}{2}-2}$ does not appear in the last integral. We can continue this procedure if needed. After $[\frac{n}{2}]$ steps we get

$$\zeta_{\alpha,n} r^{-\frac{n}{2}} + r^{-\frac{n}{2}} \int_0^1 g_{[\frac{n}{2}]}(s) J_0(rs) ds = \zeta_{\alpha,n} r^{-\frac{n}{2}} + O(r^{-\frac{n+2}{2}})$$

when n is even, and the integral

$$r^{-\frac{n-1}{2}} \int_0^1 g_{[\frac{n}{2}]}(s) s^{\frac{1}{2}} J_{\frac{1}{2}}(rs) ds = \zeta_{\alpha,n} r^{-\frac{n}{2}} + O(r^{-\frac{n+2}{2}})$$

when n is odd.

To estimate the sum we need the following lemma.

Lemma 2. For $r \geq 1$, $\beta > -\frac{1}{2}$, $\mu > -1$ and each positive integer p

$$\int_0^1 (1-s)^\mu s^{\beta+1} J_\beta(rs) ds = \sum_{j=1}^p \alpha_j^\mu r^{-(\mu+j)} J_{\beta+\mu+j}(r) + O\left(r^{-\mu-p-\frac{3}{2}}\right),$$

where α_j^μ are some numbers depending only on p and μ , and $\alpha_1^\mu = \Gamma(\mu+1)$, $\alpha_2^\mu = \mu\Gamma(\mu+2)$.

Proof. We have

$$\begin{aligned} (1-s)^\mu &= 2^{-\mu}(1-s^2)^\mu + [(1-s)^\mu - 2^{-\mu}(1-s^2)^\mu] \\ &= 2^{-\mu}(1-s^2)^\mu + 2^{-\mu}(1-s)^{\mu+1} \frac{2^\mu - (1+s)^\mu}{1-s}. \end{aligned}$$

Continuing this process of chipping off the binomials $(1-s^2)^{\mu+j-1}$ for $j = 1, \dots, p$, use the formula (see e.g., [SW], Ch.4, Lemma 4.13):

$$J_{\beta+\mu+1}(r) = \frac{r^{\mu+j}}{2^{\mu+j-1}\Gamma(\mu+j)} \int_0^1 J_\beta(rs) s^{\beta+1} (1-s^2)^{\mu+j-1} ds.$$

The remainder term is estimated as above, for the second integral for $q_\alpha(r)$, by integrating by parts $\mu^* + p + 1$ times. Estimates are better in this case since s here is in rather high power. The lemma is proved. \square

To finish the proof of Lemma 1 it remains to apply Lemma 2, with $\beta = \frac{n}{2}$, $\mu = \alpha - 1 + j$, $p = 1$, to the integrals in the sum for $q_\alpha(r)$. The proof of Lemma 1 is complete. \square

Remark 2. Sometimes the estimate $q_\alpha(r) = O\left(r^{-\alpha-\frac{1}{2}}\right)$ will be enough.

The following lemma is due to Trigub (see [T4], Lemma 2). Since that edition is difficult of access, we give the proof here. The lemma deals with the functions

$$i(\mu, \lambda, r) = \int_0^1 t^\mu J_\lambda(rt) dt$$

where $\mu + \lambda > -1$

Lemma 3.

- 1) $i(\mu, \lambda, r) = \frac{1}{r} J_{\lambda+1}(r) + \frac{\lambda+1-\mu}{r} i(\mu-1, \lambda+1, r).$
- 2) *The function $i(\mu, \lambda, r)$ is $O(r^\lambda)$ for small r , and when $r \rightarrow \infty$ it behaves as $O(r^{-\frac{3}{2}})$ or $O(r^{-1-\mu})$ for $\mu > \frac{1}{2}$ and $\mu \leq \frac{1}{2}$, respectively.*

Proof. To prove 1), integrate by parts using (7). We have

$$\begin{aligned} i(\mu, \lambda, r) &= \frac{1}{r} \int_0^1 t^{\mu-\lambda-1} d[t^{\lambda+1} J_{\lambda+1}(rt)] \\ &= \frac{1}{r} J_{\lambda+1}(r) + \frac{\lambda+1-\mu}{r} \int_0^1 t^{\mu-1} J_{\lambda+1}(rt) dt. \end{aligned}$$

The first assertion in 2) immediately follows from (9) applied to J_λ and the usual estimate of the integral. Now let $r \geq 1$. After a linear change of variables

$$i(\mu, \lambda, r) = r^{-1-\mu} \int_0^r t^\mu J_\lambda(t) dt.$$

Decompose the integral. It is bounded when $t \in [0, 1]$. When $t \in [1, r]$ use (8). If $\mu > \frac{1}{2}$, then we obtain after integrating by parts

$$\begin{aligned} \int_1^r t^\mu J_\lambda(t) dt &= \int_1^r t^\mu \left[\sqrt{\frac{2}{\pi}} \frac{\cos\left(t - \frac{\pi\lambda}{2} - \frac{\pi}{4}\right)}{\sqrt{t}} + O\left(t^{-\frac{3}{2}}\right) \right] dt \\ &= \sqrt{\frac{2}{\pi}} r^{\mu-\frac{1}{2}} \sin\left(r - \frac{\pi\lambda}{2} - \frac{\pi}{4}\right) + O(1) + \int_1^r O\left(t^{\mu-\frac{3}{2}}\right) dt = O\left(r^{\mu-\frac{1}{2}}\right), \end{aligned}$$

and

$$i(\mu, \lambda, r) = O\left(r^{-1-\mu} r^{\mu-\frac{1}{2}}\right) = O\left(r^{-\frac{3}{2}}\right).$$

If $\mu \leq \frac{1}{2}$, then the same computations show that the integral $\int_1^r t^\mu J_\lambda(t) dt$ is bounded with respect to r and

$$i(\mu, \lambda, r) = O\left(r^{-1-\mu}\right).$$

The lemma is proved. \square

3. PROOF OF THEOREM 1.

3.1. Let us start with proving (6). Due to [M], p. 140, for each $\varphi \in S$

$$\begin{aligned} (\tilde{f}, \varphi) &= \int_{\mathbb{R}^n} \tilde{f}(x) \varphi(x) dx \\ &= \lim_{A \rightarrow \infty} \int \varphi(x) \left[(2\pi)^{\frac{n}{2}} |x|^{1-\frac{n}{2}} \int_0^A f_0(t) t^{\frac{n}{2}} J_{\frac{n}{2}-1}(|x|t) dt \right] dx. \end{aligned}$$

This equality may be rewritten as follows:

$$(\tilde{f}, \varphi) = (2\pi)^{\frac{n}{2}} \lim_{A \rightarrow \infty} \int_0^\infty \varphi_0(r) r^{\frac{n}{2}} dr \int_0^A f_0(t) t^{\frac{n}{2}} J_{\frac{n}{2}-1}(rt) dt.$$

To prove (6) we will use the following procedure: integration by parts so many times that the right-hand side of (6) is obtained, and then justification of the passage to limit under the integral sign. Thus, for $[\alpha] \neq 0$ integrate in t by parts $[\alpha]$ times. We obtain

$$\begin{aligned} (10) \quad & \int_0^\infty \varphi_0(r) r^{\frac{n}{2}} dr \int_0^A f_0(t) t^{\frac{n}{2}} J_{\frac{n}{2}-1}(rt) dt \\ &= \int_0^\infty \varphi_0(r) r^{\frac{n}{2}} \left\{ \frac{(-1)^{[\alpha]}}{([\alpha]-1)!} \int_0^A f_0^{([\alpha])}(t) t^{[\alpha]+\frac{n}{2}} dt \int_0^1 (1-s)^{[\alpha]-1} s^{\frac{n}{2}} J_{\frac{n}{2}-1}(rts) ds \right. \\ & \quad \left. + \sum_{p=0}^{[\alpha]-1} \frac{(-1)^p}{p!} f_0^{(p)}(t) t^{p+\frac{n}{2}+1} \int_0^1 (1-s)^p s^{\frac{n}{2}} J_{\frac{n}{2}-1}(rts) ds \right\} dr. \end{aligned}$$

To show that integrated terms vanish when $t = 0$ apply (9) to $J_{\frac{n}{2}-1}(rts)$. Let us demonstrate that

$$\lim_{t \rightarrow 0+} f_0^{(p)}(t) t^{p+n} = 0 \quad \text{for} \quad p = 0, 1, \dots, [\alpha] - 1 \quad \text{and} \quad \alpha \geq 1.$$

For α fractional we have

$$f_0^{([\alpha])}(t) = \frac{1}{\Gamma(\alpha - [\alpha])} \int_t^\infty (s-t)^{\alpha-[\alpha]-1} f_0^{(\alpha)}(s) ds.$$

This formula due to Cossar [Co] may be found also in [Tr], Lemma 3.10. It is easy to see now that

$$f_0^{(p)}(t) = \frac{(-1)^{[\alpha]-p}}{\Gamma(\alpha - p)} \int_t^\infty (s-t)^{\alpha-p-1} f_0^{(\alpha)}(s) ds.$$

It is clear that for α integer the same formula holds. Integrate by parts as follows

$$\begin{aligned} \int_t^\infty (s-t)^{\alpha-p-1} f_0^{(\alpha)}(s) ds &= -\frac{1}{\alpha-p} (s-t)^{\alpha-p} f_0^{(\alpha)}(s) ds \Big|_t^\infty \\ &+ \frac{1}{\alpha-p} \int_t^\infty (s-t)^{\alpha-p} f_0^{(\alpha+1)}(s) ds \\ &= \frac{1}{\alpha-p} \int_t^\infty (s-t)^{\alpha-p} \frac{1}{s^{\frac{n-1}{2}}} dF_\alpha(s) \\ &- \frac{1}{\alpha-p} \frac{n-1}{2} \int_t^\infty (s-t)^{\alpha-p} \frac{1}{s} f_0^{(\alpha)}(s) ds. \end{aligned}$$

Here and below (3) is taken into account for integrated terms. Split the last integral into two one of which is the same as that on the left-hand side. Taking into account the corresponding coefficient we get

$$\begin{aligned} \left(1 + \frac{1}{\alpha - p} \frac{n-1}{2}\right) \int_t^\infty (s-t)^{\alpha-p-1} f_0^{(\alpha)}(s) ds &= \frac{1}{\alpha - p} \int_t^\infty (s-t)^{\alpha-p} \frac{1}{s^{\frac{n-1}{2}}} dF_\alpha(s) \\ &+ \frac{t}{\alpha - p} \frac{n-1}{2} \int_t^\infty (s-t)^{\alpha-p-1} \frac{1}{s} f_0^{(\alpha)}(s) ds. \end{aligned}$$

The last integral is equal to

$$\int_t^\infty (s-t)^{\alpha-p-1} \frac{1}{s^{\frac{n+1}{2}}} F_\alpha(s) ds.$$

Integrate by parts again and obtain

$$\begin{aligned} \int_t^\infty (s-t)^{\alpha-p-1} \frac{1}{s^{\frac{n+1}{2}}} F_\alpha(s) ds &= -F_\alpha(s) \int_s^\infty (u-t)^{\alpha-p-1} \frac{1}{u^{\frac{n+1}{2}}} du \Big|_t^\infty \\ &+ \int_t^\infty dF_\alpha(s) \int_s^\infty (u-t)^{\alpha-p-1} \frac{1}{u^{\frac{n+1}{2}}} du \\ &= F_\alpha(t) \int_t^\infty (u-t)^{\alpha-p-1} \frac{1}{u^{\frac{n+1}{2}}} du \\ &+ \int_t^\infty dF_\alpha(s) \int_s^\infty (u-t)^{\alpha-p-1} \frac{1}{u^{\frac{n+1}{2}}} du. \end{aligned}$$

Let us consider the integral

$$\int_s^\infty (u-t)^{\alpha-p-1} \frac{1}{u^{\frac{n+1}{2}}} du = \int_s^\infty u^{-\frac{n+1}{2} + \alpha - p - 1} du.$$

It can be obviously estimated by $t^{-\frac{n+1}{2} + \alpha - p}$ and now condition (4) completes the proof.

These calculations may seem somewhat superfluous. Indeed, for most cases an easier way gives the result needed but for $p = 0$ when $\alpha = \frac{n-1}{2}$ one has to be more careful.

For $t = A$, taking into account that $A \rightarrow \infty$, it suffices, in view of (2), to establish the uniform boundedness in t of the integrals

$$B_p(t) = \int_0^\infty \varphi_0(r) r^{\frac{n}{2}} dr t^{\frac{n}{2}+1} \int_0^1 (1-s)^p s^{\frac{n}{2}} J_{\frac{n}{2}-1}(rts) ds.$$

Integrate by parts in the outer integral $m = \left[\frac{n-1}{2}\right]$ times, using (7) so that the order of the Bessel function decreases. Integrated terms vanish since $\varphi_0 \in S$. We have (defining by ψ here and below some function from S)

$$B_p(t) = \int_0^\infty \psi(r) r^{\frac{n}{2}-m} t^{\frac{n}{2}-m+1} dr \int_0^1 (1-s)^p s^{\frac{n}{2}-m} J_{\frac{n}{2}-m-1}(rts) ds.$$

For n odd

$$B_p(t) = t \int_0^\infty \psi(r) dr \int_0^1 (1-s)^p \cos rts ds.$$

For $p = 0$ we have the Fourier integral formula (see [Bo], §9) : $\lim_{t \rightarrow \infty} B_0(t) = \pi^{-1} \psi(0)$. For $p \geq 1$

$$\left| t \int_1^\infty \psi(r) dr \int_0^1 (1-s)^p \cos rts ds \right| \leq \int_1^\infty |\psi(r)| \frac{dr}{r} < \infty$$

and

$$\begin{aligned} t \int_0^1 \psi(r) dr \int_0^1 (1-s)^p \cos rts ds &= t \int_0^1 \psi(0) dr \int_0^1 (1-s)^p \cos rts ds \\ &+ t \int_0^1 [\psi(r) - \psi(0)] dr \int_0^1 (1-s)^p \cos rts ds. \end{aligned}$$

The first integral on the right-hand side is equal to

$$\psi(0) \int_0^1 (1-s)^p s^{-1} \sin ts ds,$$

and its finiteness is well-known. For the second integral the estimate

$$\int_0^1 |\psi(r) - \psi(0)| \frac{dr}{r} < \infty$$

follows, as above. Now for n even

$$\begin{aligned} B_p(t) &= \int_0^\infty \psi(r) r t^2 dr \int_0^1 (1-s)^p s J_0(rts) ds \\ &= - \int_0^\infty \psi'(r) \left[r t \int_0^1 (1-s)^p s J_1(rts) ds \right] dr, \end{aligned}$$

and integration by parts due to (7) yields the boundedness of $B_p(t)$ immediately. Therefore, the integrated terms in (10) vanish as $A \rightarrow \infty$. For α integer, the element of integration coincides with that indicated in (6).

When α is fractional, apply to the first summand in the curly brackets on the right-hand side of (10) the following formula of fractional integration by parts (see [BE2], p. 182, or [SKM], (2.20))

$$\int_0^\infty f_1(t) R_\gamma(f_2; t) dt = \int_0^\infty W_\gamma(f_1; t) f_2(t) dt.$$

In our case $\gamma = 1 - \alpha + [\alpha]$,

$$f_1(t) = \begin{cases} f^{([\alpha])}(t), & t \leq A, \\ 0, & t > A, \end{cases}$$

and $f_2(t) = \frac{\Gamma([\alpha])}{\Gamma(\alpha)} \frac{d}{dt} [t^{\alpha+\frac{n}{2}} Q_\alpha(rt)]$. Indeed,

$$\begin{aligned} R_{\alpha-[\alpha]}(s^{[\alpha]+\frac{n}{2}} Q_{[\alpha]}(rs); t) &= \Gamma([\alpha]) R_{\alpha-[\alpha]}(R_{[\alpha]}(s^{\frac{n}{2}} J_{\frac{n}{2}-1}(rs); \cdot); t) \\ &= \Gamma([\alpha]) R_\alpha(s^{\frac{n}{2}} J_{\frac{n}{2}-1}(rs); t) = \frac{\Gamma([\alpha])}{\Gamma(\alpha)} t^{\alpha+\frac{n}{2}} Q_\alpha(rt). \end{aligned}$$

Therefore f_2 is the Riemann-Liouville derivative of order $1 - \alpha + [\alpha]$ of the function $t^{[\alpha]+\frac{n}{2}} Q_{[\alpha]}(rt)$ and the Riemann-Liouville integral of order $1 - \alpha + [\alpha]$ of f_2 is exactly $t^{[\alpha]+\frac{n}{2}} Q_{[\alpha]}(rt)$.

Apply the usual integration by parts to the right-hand side of the formula of the fractional integration by parts. This gives the following equality true for all α :

$$\begin{aligned} & \frac{1}{([\alpha] - 1)!} \int_0^A f_0^{([\alpha])}(t) t^{[\alpha]+\frac{n}{2}} Q_{[\alpha]}(rt) dt \\ &= \frac{1}{\Gamma(\alpha)} \int_0^A \frac{d}{dt} [t^{\alpha+\frac{n}{2}} Q_\alpha(rt)] dt \frac{1}{\Gamma(1 - \alpha + [\alpha])} \int_t^A f^{([\alpha])}(s) (s - t)^{[\alpha]-\alpha} ds \\ &= \frac{1}{\Gamma(\alpha) \Gamma(1 - \alpha + [\alpha])} \left[\int_t^A (s - t)^{-\alpha+[\alpha]} f_0^{([\alpha])}(s) ds \right] t^{\alpha+\frac{n}{2}} Q_\alpha(rt) \Big|_0^A \\ &+ \frac{[\alpha] - \alpha}{\Gamma(\alpha) \Gamma(1 - \alpha + [\alpha])} \int_0^A t^{\alpha+\frac{n}{2}} Q_\alpha(rt) dt \int_A^\infty (s - t)^{-\alpha+[\alpha]-1} f_0^{([\alpha])}(s) ds \\ &- \frac{1}{\Gamma(\alpha)} \int_0^A F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt. \end{aligned}$$

It must be shown again that the two first values on the right-hand side vanish as $A \rightarrow \infty$. For t large enough

$$\begin{aligned} & t^{\frac{n}{2}+\alpha} \left| \int_t^A (s - t)^{-\alpha+[\alpha]} f^{([\alpha])}(s) ds \right| \\ & \leq t^{\frac{n}{2}+\alpha-[\alpha]} \int_t^A (s - t)^{-\alpha+[\alpha]} |s^{[\alpha]} f^{([\alpha])}(s)| ds \end{aligned}$$

and because of (2) the right-hand side equals to zero for $t = A$. Consider now

$$t^{\frac{n}{2}+\alpha} Q_\alpha(rt) \int_t^A (s - t)^{-\alpha+[\alpha]} f^{([\alpha])}(s) ds$$

as $t \rightarrow 0$. We will use the inequality

$$Q_\alpha(rt) \leq (Cr^{\frac{n}{2}-1})t^{\frac{n}{2}-1}$$

that immediately follows from (9). As above

$$\begin{aligned} & \left| t^{\frac{n}{2}+\alpha} Q_\alpha(rt) \int_1^A (s-t)^{-\alpha+[\alpha]} f^{([\alpha])}(s) ds \right| \\ & \leq C(A-t)^{1-\alpha+[\alpha]} t^{n-1+\alpha-[\alpha]} \sup_{[1,A]} |s^{[\alpha]} f^{([\alpha])}(s)| \end{aligned}$$

and the right-hand side tends to zero as $t \rightarrow 0$. Estimate now

$$t^{\frac{n}{2}+\alpha} Q_\alpha(rt) \int_t^1 (s-t)^{-\alpha+[\alpha]} f^{([\alpha])}(s) ds.$$

We wish to show that

$$\lim_{t \rightarrow 0+} t^{\alpha+n-1} \int_t^1 (s-t)^{[\alpha]-\alpha} f_0^{([\alpha])}(s) ds = 0.$$

From the definition of F_α we obtain

$$\frac{d}{dt} \int_t^\infty (s-t)^{[\alpha]-\alpha} f_0^{([\alpha])}(s) ds = O(t^{-\frac{n-1}{2}}).$$

Since

$$\frac{d}{dt} \int_1^\infty (s-t)^{[\alpha]-\alpha} f_0^{([\alpha])}(s) ds = (\alpha - [\alpha]) \int_1^\infty (s-t)^{[\alpha]-\alpha-1} f_0^{([\alpha])}(s) ds$$

is bounded, we have

$$\frac{d}{dt} \int_t^1 (s-t)^{[\alpha]-\alpha} f_0^{([\alpha])}(s) ds = O(t^{-\frac{n-1}{2}})$$

and thus

$$\int_t^1 (s-t)^{[\alpha]-\alpha} f_0^{([\alpha])}(s) ds = O(t^{-\frac{n-3}{2}}).$$

This estimate proves the statement

Further, integrating, in r , by parts $[\frac{n-1}{2}]$ times using (7) with “-” in the left-hand side, and then once more using (7) with “+” in the left-hand side, we obtain

$$\begin{aligned} & \int_0^\infty \varphi_0(r) r^{\frac{n}{2}} dr \int_0^A \left\{ \int_A^\infty (s-t)^{-\alpha+[\alpha]-1} f_0^{([\alpha])}(s) ds \right\} t^{\frac{n}{2}+\alpha} Q_\alpha(rt) dt \\ &= \int_0^\infty \psi(r) r^{\frac{n}{2}-[\frac{n+1}{2}]+1} dr \int_0^A \left\{ \int_A^\infty (s-t)^{-\alpha+[\alpha]-1} f_0^{([\alpha])}(s) ds \right\} \\ & \quad \times t^{\frac{n}{2}-[\frac{n+1}{2}]+\alpha} dt \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-[\frac{n+1}{2}]} J_{\frac{n}{2}+1-[\frac{n+1}{2}]}(rtu) du. \end{aligned}$$

In view of (2) we have $\sup_{s \in [A, \infty)} \left| s^{[\alpha]} f_0^{([\alpha])}(s) \right| \longrightarrow O$ as $A \rightarrow \infty$. Besides that

$$\begin{aligned} & \int_0^A t^{\alpha-[\alpha]-1} dt \int_A^\infty (s-t)^{-\alpha+[\alpha]-1} ds \\ &= \frac{1}{\alpha-[\alpha]} \int_0^1 t^{\alpha-[\alpha]-1} (1-t)^{-\alpha+[\alpha]} dt \leq C. \end{aligned}$$

For n odd and each α , since $J_{\frac{1}{2}}(rtu) = \sqrt{\frac{2}{\pi rtu}} \sin rtu$, we obtain

$$\left| (rt)^{\frac{n}{2}-[\frac{n+1}{2}]+1} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-[\frac{n+1}{2}]} J_{\frac{n}{2}+1-[\frac{n+1}{2}]}(rtu) dr \right| \leq C.$$

For n even Lemma 1, with $n = 2$, yields such an estimate only for $\alpha \geq \frac{1}{2}$. It remains to consider

$$\int_0^\infty \psi(r) r dr \int_0^A \left\{ \int_A^\infty (s-t)^{-\alpha-1} f_0^{([\alpha])}(s) ds \right\} t^\alpha dt \int_0^1 (1-u)^{\alpha-1} J_1(rt) du$$

for $0 < \alpha < \frac{1}{2}$. Again, applying Lemma 1 with $n = 2$, we reduce the problem to finiteness of

$$\int_1^\infty \psi(r) (rt)^{\frac{1}{2}-\alpha} \sin rt dr$$

when $rt > 1$, since all the remainder terms are estimated as above after using (7). Let us integrate by parts. All the integrated terms are bounded, and it remains to estimate

$$\int_1^\infty \psi'(r) r (rt)^{-\frac{1}{2}-\alpha} \cos rt dr + \left(\frac{1}{2} - \alpha \right) \int_1^\infty \psi(r) (rt)^{-\frac{1}{2}-\alpha} \cos rt dr.$$

Both integrals are finite, and we get finally

$$(\tilde{f}, \varphi) = \frac{(2\pi)^{\frac{n}{2}} (-1)^{\alpha^*+1}}{\Gamma(\alpha)} \lim_{A \rightarrow \infty} \int_0^\infty \varphi_0(r) r^{\frac{n}{2}} dr \int_0^A F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt.$$

It remains to justify the passage to the limit under the integral sign. Due to the Lebesgue dominated theorem, it is possible if the element of integration is dominated by an integrable function independent of A . To prove this, consider two integrals: over $r \in [0, 1]$ and over $r \in [1, \infty)$, respectively. In view of (3) and (4), one can treat F_α as a function which is monotone decreasing and vanishing at infinity. Let $r \in [1, \infty)$. Consider two integrals in t : over $[0, 1]$ and $[1, A]$, respectively. The first one is bounded, and the majorant is simply $C |\varphi_0(r) r^{\frac{n}{2}}|$. Apply the second mean value theorem of integral calculus to the integral over $[1, A]$. Using (7) and Lemma 1, we obtain ($\xi \leq A$):

$$\begin{aligned} & \left| \varphi_0(r) r^{\frac{n}{2}} \int_1^A F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt \right| = \left| \varphi_0(r) r^{\frac{n}{2}} F_\alpha(1) \int_1^\xi t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt \right| \\ &= \left| F_\alpha(1) \varphi_0(r) r^{\frac{n}{2}-1} \left\{ t^{\alpha+\frac{1}{2}} q_\alpha(rt) \Big|_1^\xi + \left(\frac{n-1}{2} - \alpha \right) \int_1^\xi t^{\alpha-\frac{1}{2}} q_\alpha(rt) dt \right\} \right| \\ &\leq \left(\frac{n-1}{2} - \alpha \right) \left| F_\alpha(1) \varphi_0(r) r^{\frac{n-3}{2}-\alpha} \int_1^\xi \frac{1}{t} \cos(rt + \mu) dt \right| + C |F_\alpha(1) \varphi_0(r)| \frac{1}{r}, \end{aligned}$$

and the last value is integrable over $[1, \infty)$. At last

$$\begin{aligned} & \lim_{A \rightarrow \infty} \int_0^1 \varphi_0(r) r^{\frac{n}{2}} dr \int_0^A F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt \\ &= \lim_{A \rightarrow \infty} \int_0^1 \{[\varphi_0(r) - \varphi_0(0)] + \varphi_0(0)\} r^{\frac{n}{2}} dt \int_0^A F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt. \end{aligned}$$

Since $\frac{1}{r} [\varphi_0(r) - \varphi_0(0)]$ is integrable, the part corresponding to this function may be estimated completely like in the case $r \in [1, \infty)$. Further,

$$\lim_{A \rightarrow \infty} \int_0^1 r^{\frac{n}{2}} \int_0^A F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt = \lim_{A \rightarrow \infty} \int_0^A F_\alpha(t) t^{\alpha-\frac{1}{2}} q_\alpha(t) dt.$$

Since F_α is bounded and monotone, integration by parts and estimates like in Lemma 1 yield the convergence of this integral in improper sense.

In fact, we have proved the uniform convergence of the integral (6) when $|x| \geq r_0 > 0$.

3.2. Let us show now that $\hat{f}(x) \rightarrow 0$ as $r = |x| \rightarrow \infty$. Using the second mean value theorem, we obtain for some $A'' \leq A'$:

$$r^{1-\frac{n}{2}} \int_A^{A'} F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt = F_\alpha(A) r^{-\frac{n}{2}} \left[t^{\alpha+\frac{1}{2}} q_\alpha(rt) \right]_A^{A'}.$$

In view of Lemma 1, we have for every $A \geq 1$ as $A' \rightarrow \infty$

$$\left| r^{1-\frac{n}{2}} \int_A^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt \right| \leq C |F_\alpha(A)| r^{-\frac{n+1}{2}-\alpha}.$$

The right-hand side tends to zero as $r \rightarrow \infty$. Further,

$$\begin{aligned}
& r^{1-\frac{n}{2}} \int_0^A F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(rt) dt \\
&= F_\alpha(t) r^{-\frac{n}{2}} q_\alpha(rt) \Big|_0^A - r^{-\frac{n}{2}} \int_0^A t^{\alpha+\frac{1}{2}} q_\alpha(rt) dF_\alpha(t) \\
&+ \left(\frac{n-1}{2} - \alpha \right) r^{-\frac{n}{2}} \int_0^A F_\alpha(t) t^{\alpha-\frac{1}{2}} q_\alpha(rt) dt = O(r^{-\frac{n}{2}})
\end{aligned}$$

by Lemma 1 and (4).

3.3. Let us show the continuity of $\hat{f}(x)$ for $|x| > 0$. Let $[r_0, r_1]$ be an interval of uniform convergence of the integral in (6), and $|x| \in [r_0, r_1]$. Then the functions

$$\hat{f}_k(x) = \frac{(2\pi)^{\frac{n}{2}} (-1)^{\alpha^*+1}}{\Gamma(\alpha)} |x|^{1-\frac{n}{2}} \int_0^k F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(|x|t) dt$$

are continuous for each $k = 1, 2, \dots$, and converge uniformly to $\hat{f}(x)$ as $k \rightarrow \infty$. Hence, $\hat{f}(x)$ is continuous for these x as well.

3.4. Let us prove now the inverse formula. Applying the Cauchy-Poisson formula (see e.g., [Bo], Th. 56, or [SW], Ch.4, Th.3.3), we have

$$\begin{aligned}
(11) \quad & (2\pi)^{-n} \int_{|u| \leq A} \left(1 - \frac{|u|^2}{A^2} \right)^{\frac{n-1}{2}-\alpha} \hat{f}(u) e^{ixu} du \\
&= \frac{(-1)^{\alpha^*+1}}{\Gamma(\alpha)} r^{1-\frac{n}{2}} \int_0^A \left(1 - \frac{s^2}{A^2} \right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) ds \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(st) dt \\
&= \frac{(-1)^{\alpha^*+1}}{\Gamma(\alpha)} r^{1-\frac{n}{2}} \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} dt \int_0^A \left(1 - \frac{s^2}{A^2} \right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds.
\end{aligned}$$

This change of integrals must be justified. Let $0 < \delta < A$. The uniform convergence of the integral in t for $s \geq \delta$ yields

$$\begin{aligned}
(12) \quad & \int_\delta^A \left(1 - \frac{s^2}{A^2} \right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) ds \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} Q_\alpha(st) dt \\
&= \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} dt \int_0^A \left(1 - \frac{s^2}{A^2} \right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds \\
&- \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} dt \int_0^\delta \left(1 - \frac{s^2}{A^2} \right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds.
\end{aligned}$$

It suffices to show that the last integral tends to zero as $\delta \rightarrow 0$. Take $\varepsilon > 0$ and let M be large enough to provide $|F_\alpha(M)| < \varepsilon$, by (3). The second mean value theorem yields after integrating by parts:

$$\begin{aligned}
(13) \quad & \int_M^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} dt \int_0^\delta \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds \\
&= F_\alpha(M) \int_M^{M'} t^{\alpha+\frac{1}{2}} dt \int_0^\delta \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds \\
&= F_\alpha(M) \left[t^{\alpha+\frac{1}{2}} \int_0^\delta \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} J_{\frac{n}{2}-1}(rs) q_\alpha(st) ds \right]_M^{M'} \\
&\quad - F_\alpha(M) \left(\alpha - \frac{n-1}{2} \right) \int_M^{M'} t^{\alpha-\frac{1}{2}} dt \int_0^\delta \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} J_{\frac{n}{2}-1}(rs) q_\alpha(st) ds.
\end{aligned}$$

Estimate integrated terms in (13). The uniform boundedness, in t and δ , of the value in brackets should be shown. Since $\zeta_{\alpha,n} \int_0^\delta \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s^{-\alpha-\frac{1}{2}} J_{\frac{n}{2}-1}(rs) ds$ does not depend on M and M' , such values, taking twice with opposite signs, are cancelled. The rest, in view of (9), does not exceed, for $s \in [0, \frac{1}{t}]$, the quantity

$$C t^{\alpha+\frac{1}{2}} \int_0^{\frac{1}{t}} |J_{\frac{n}{2}-1}(rs)| ds \leq C r^{\frac{n}{2}-1} t^{\alpha-\frac{n-1}{2}},$$

and for $t\delta > 1$ and $s \in [\frac{1}{t}, \delta]$, in view of Lemma 1 and (9), is

$$\Gamma(\alpha) t^{\frac{1}{2}} \int_{\frac{1}{t}}^\delta s^{-\alpha} \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}+\alpha}(st) ds + O\left(\frac{1}{t} \int_{\frac{1}{t}}^\delta s^{\frac{n-5}{2}-\alpha} ds\right).$$

The remainder term does not exceed $C \left\{ t^{\alpha-\frac{n-1}{2}} + (t\delta)^{-1} \delta^{\frac{n-1}{2}-\alpha} \right\}$, which is bounded. Apply (9) to $J_{\frac{n}{2}-1}$ and (8) to $J_{\frac{n}{2}+\alpha}$ in the main term. We obtain for $\alpha < \frac{n-1}{2}$

$$\int_{\frac{1}{t}}^\delta \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s^{-\alpha+\frac{n}{2}-\frac{3}{2}} ds \leq C \left(\delta^{\frac{n-1}{2}-\alpha} + t^{\alpha-\frac{n-1}{2}} \right) \leq C.$$

For $\alpha = \frac{n-1}{2}$, integrate by parts as follows:

$$\begin{aligned}
& t^{\frac{1}{2}} \int_{\frac{1}{t}}^\delta s^{-\frac{n-1}{2}} J_{\frac{n}{2}-1}(rs) J_{n-\frac{1}{2}}(st) ds \\
&= -t^{-\frac{1}{2}} s^{-\frac{n-1}{2}} J_{\frac{n}{2}-1}(rs) J_{n+\frac{1}{2}}(st) \Big|_{\frac{1}{t}}^\delta + r t^{-\frac{1}{2}} \int_{\frac{1}{t}}^\delta s^{-\frac{n-1}{2}} J_{\frac{n}{2}-2}(rs) J_{n+\frac{1}{2}}(st) ds,
\end{aligned}$$

and now the proof is continued as for the remainder term.

Let us estimate the last integral in (13). This makes sense only for $\alpha < \frac{n-1}{2}$. Use again Lemma 1. We have to estimate

$$\int_M^{M'} t^{-\frac{1}{2}} dt \int_{\frac{1}{t}}^{\delta} \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s^{-\alpha} J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}+\alpha}(ts) ds,$$

since the rest is estimated analogously. Let us change the order of integration. Without loss of generality, one can take $\delta < \frac{1}{M}$. The following should be estimated:

$$\begin{aligned} & \int_{\frac{1}{M'}}^{\delta} \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s^{-\alpha} J_{\frac{n}{2}-1}(rs) ds \int_{\frac{1}{s}}^{M'} t^{-\frac{1}{2}} J_{\frac{n}{2}+\alpha}(st) dt \\ & + \int_{\delta}^{\frac{1}{M}} \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s^{-\alpha} J_{\frac{n}{2}-1}(rs) ds \int_M^{\frac{1}{s}} t^{-\frac{1}{2}} J_{\frac{n}{2}+\alpha}(st) dt. \end{aligned}$$

Apply (8) to $J_{\frac{n}{2}+\alpha}$. For the remainder term, after applying (9) to $J_{\frac{n}{2}-1}$, we obtain

$$\int_{\frac{1}{M'}}^{\delta} s^{\frac{n-5}{2}-\alpha} ds \int_{\frac{1}{s}}^{M'} \frac{dt}{t^2} + \int_{\delta}^{\frac{1}{M}} s^{\frac{n-5}{2}-\alpha} ds \int_M^{\frac{1}{s}} \frac{dt}{t^2}$$

which is obviously bounded. For the main term, the integral in t is of the form

$$\int t^{-1} \cos(st + \mu) dt.$$

After integrating by parts the estimates coincide with those for the remainder term. Hence, the value (13) is small. Choosing δ so small that

$$\left| \int_0^M F_{\alpha}(t) t^{\alpha+\frac{1}{2}} dt \int_0^{\delta} \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) Q_{\alpha}(st) ds \right| < \varepsilon,$$

we obtain that the last integral in (12) tends to zero as $\delta \rightarrow 0$. Returning to (11) we have

$$\begin{aligned} & \int_0^A \left(1 - \frac{s^2}{A^2}\right)^{\frac{n-1}{2}-\alpha} s J_{\frac{n}{2}-1}(rs) Q_{\alpha}(st) ds = \int_0^A s J_{\frac{n}{2}-1}(rs) Q_{\alpha}(st) ds \\ & - \int_0^A s J_{\frac{n}{2}-1}(rs) Q_{\alpha}(st) ds \, 2 \left(\frac{n-1}{2} - \alpha \right) \int_{\frac{s}{A}}^{\frac{s}{A}} u(1-u^2)^{\frac{n-3}{2}-\alpha} du. \end{aligned}$$

Let us begin with the second integral. We have

$$\begin{aligned}
& \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} dt \int_0^A s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds \int_0^{\frac{s}{A}} u(1-u^2)^{\frac{n-3}{2}-\alpha} du \\
&= \int_0^1 u(1-u^2)^{\frac{n-3}{2}-\alpha} du \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} dt \int_{Au}^A s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds \\
&= \int_0^1 u(1-u^2)^{\frac{n-3}{2}-\alpha} du \left\{ F_\alpha(t) t^{\alpha+\frac{1}{2}} \int_{Au}^A J_{\frac{n}{2}-1}(rs) q_\alpha(st) ds \right\}_0^\infty \\
&\quad - \int_0^\infty dF_\alpha(t) t^{\alpha+\frac{1}{2}} \int_{Au}^A J_{\frac{n}{2}-1}(rs) q_\alpha(st) ds \\
&\quad + \left(\frac{n-1}{2} - \alpha \right) \int_0^\infty F_\alpha(t) t^{\alpha-\frac{1}{2}} dt \int_{Au}^A J_{\frac{n}{2}-1}(rs) q_\alpha(st) ds \Bigg\}.
\end{aligned}$$

The integrated terms for $t = 0$ and $t = \infty$ vanish. Indeed, it is obvious for $t = 0$ and for $t = \infty$ follows from Lemma 1 and (3). We have to show that the right-hand side tends to zero as $A \rightarrow \infty$. Note firstly, that the estimate $|q_\alpha(st)| \leq C(st)^{-\alpha-\frac{1}{2}}$ and estimates (8) and (9) for $J_{\frac{n}{2}-1}$ yield

$$\left| t^{\alpha+\frac{1}{2}} \int_{Au}^A J_{\frac{n}{2}-1}(rs) q_\alpha(st) ds \right| \leq C \int_{Au}^A \frac{ds}{s^{1+\varepsilon}}$$

with some $\varepsilon \in (0, 1)$. This estimate combined with (4) gives proper estimates for the second summand in the curly brackets. When $t \in [0, 1]$ the estimates for the third one are similar. Use then Lemma 1 when $t \in [1, \infty)$. The remainder term is estimated as above. The main term, by the second mean value theorem, is equal to

$$\begin{aligned}
& \int_1^\infty F_\alpha(t) t^{-\frac{1}{2}} dt \int_{Au}^A s^{-\alpha} J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}+\alpha}(st) ds \\
&= F_\alpha(1) \int_1^\xi t^{-\frac{1}{2}} dt \int_{Au}^A s^{-\alpha} J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}+\alpha}(st) ds.
\end{aligned}$$

Apply (8) to $J_{\frac{n}{2}+\alpha}$. The remainder term does not need new techniques. It remains to estimate

$$\int_1^\xi \frac{dt}{t} \int_{Au}^A s^{-\alpha-\frac{1}{2}} J_{\frac{n}{2}-1}(rs) \cos(ts + \mu) ds.$$

Integration by parts in t and estimates of the integral in s , like above, finish the proof of a tendency of the limit to zero. It remains to consider

$$\begin{aligned}
& \int_0^\infty F_\alpha(t) t^{\alpha+\frac{1}{2}} dt \int_0^A s J_{\frac{n}{2}-1}(rs) Q_\alpha(st) ds \\
&= \int_0^\infty F_\alpha(t) t^{\alpha-\frac{n-1}{2}} \frac{d}{dt} \left[t^{\frac{n}{2}} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_0^A J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(uts) ds \right] dt.
\end{aligned}$$

Let us substitute integration over $[0, A]$ by integration over the difference of two sets: $[0, \infty)$ and $[A, \infty)$. Let us use the formula

$$\Gamma(\nu - \mu) \int_0^\infty J_\mu(at) J_\nu(bt) t^{\mu-\nu+1} dt = \begin{cases} 2^{\mu-\nu+1} a^\mu b^{-\nu} (b^2 - a^2)^{\nu-\mu+1}, & \text{for } b > a, \\ 0, & \text{for } b < a \end{cases}$$

which is true for $\nu > \mu > -1$ (see [BE2], p. 148). We obtain

$$\begin{aligned} & \int_0^\infty F_\alpha(t) t^{\alpha-\frac{n-1}{2}} \frac{d}{dt} \left[t^{\frac{n}{2}} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_0^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(uts) ds \right] dt \\ &= \int_r^\infty F_\alpha(t) t^{\alpha-\frac{n-1}{2}} \frac{d}{dt} \left[t^{\frac{n}{2}} \int_{\frac{r}{t}}^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} r^{\frac{n}{2}-1} (tu)^{-\frac{n}{2}} du \right] dt \\ &= r^{\frac{n}{2}-1} \int_r^\infty f_0^{(\alpha)}(t) (t-r)^{\alpha-1} dt. \end{aligned}$$

Integrating by parts and using (1)–(3), we get

$$\frac{(-1)^{\alpha^*+1}}{\Gamma(a)} \int_r^\infty f_0^{(\alpha)}(t) (t-r)^{\alpha-1} dt = f_0(r).$$

For α fractional we used also the permutability of the fractional integral and fractional derivative. If we show that

$$\int_0^\infty F_\alpha(t) t^{\alpha-\frac{n-1}{2}} \frac{d}{dt} \left[t^{\frac{n}{2}} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(uts) ds \right] dt$$

tends to zero as $A \rightarrow \infty$ the inverse formula will be proved. Integration by parts in the outer integral yields

$$\begin{aligned} & F_\alpha(t) t^{\alpha+\frac{1}{2}} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(ust) ds \Big|_0^\infty \\ & - \int_0^\infty \left[t^{\alpha+\frac{1}{2}} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(ust) ds \right] dF_\alpha(t) \\ & + \left(\frac{n-1}{2} - \alpha \right) \int_0^\infty F_\alpha(t) t^{\alpha-\frac{1}{2}} dt \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(ust) ds. \end{aligned}$$

To estimate the first two summands, in view of (4), it suffices to show that

$$\limsup \left| t^{\alpha+\frac{1}{2}} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(ust) ds \right| = 0.$$

Let us show that the change of the order of integration is legal. By the Lebesgue dominated theorem, it suffices to find an integrable majorant, independent of A and A' , for the function

$$(1-u)^{\alpha-1} u^{\frac{n}{2}-1} \int_A^{A'} J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(uts) ds.$$

Integrating by parts in the inner integral, we obtain

$$\begin{aligned} & \int_A^{A'} J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(uts) ds \\ &= \frac{1}{r} J_{\frac{n}{2}}(rs) J_{\frac{n}{2}}(uts) \Big|_A^{A'} + \frac{tu}{r} \int_A^{A'} J_{\frac{n}{2}}(rs) J_{\frac{n}{2}+1}(uts) ds. \end{aligned}$$

The integrated terms are bounded, and $(1-u)^{\alpha-1} u^{\frac{n}{2}-1}$ is integrable over $[0, 1]$. Apply (8) to the Bessel functions in the last integral. The remainder terms are bounded, and $(1-u)^{\alpha-1} u^{\frac{n-3}{2}}$ is integrable. We have for the main terms

$$\begin{aligned} & \int_A^{A'} \frac{\cos(rs - \frac{\pi n}{2} - \frac{\pi}{4}) \cos(uts - \frac{\pi n}{2} - \frac{3\pi}{4})}{s} ds \\ &= \frac{1}{2} \int_A^{A'} \frac{\cos(rs + uts - \pi n - \pi)}{s} ds - \frac{1}{2} \int_A^{A'} \frac{\sin s(r - tu)}{s} ds, \end{aligned}$$

and the boundedness of these integrals is easily obtained by integration by parts and inequality $(r + ut)^{-1} \leq r^{-1}$. Now Lemma 1 and the rough estimate $J_\nu(t) = O(t^{-\frac{1}{2}})$ yield

$$\begin{aligned} & \lim_{A \rightarrow \infty} \sup_t \left| t^{\alpha+\frac{1}{2}} \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(uts) ds \right| \\ &= \lim_{A \rightarrow \infty} \sup_t \left| t^{\alpha+\frac{1}{2}} \int_A^\infty J_{\frac{n}{2}-1}(rs) q_\alpha(ts) ds \right| \leq C \lim_{A \rightarrow \infty} r^{-\frac{1}{2}} \int_A^\infty \frac{ds}{s^{1+\alpha}} = 0. \end{aligned}$$

It remains to estimate

$$\begin{aligned} & \int_0^\infty F_\alpha(t) t^{\alpha-\frac{1}{2}} dt \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(uts) ds \\ &= \int_0^\infty F_\alpha(t) t^{\alpha-\frac{1}{2}} dt \int_A^\infty J_{\frac{n}{2}-1}(rs) q_\alpha(ts) ds. \end{aligned}$$

For $t \in [\frac{r}{2}, \infty]$ such estimates are already made above. Let $t \in [0, \frac{r}{2}]$. Consider

$$\int_0^{\frac{r}{2}} F_\alpha(t) t^{\alpha-\frac{1}{2}} dt \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty J_{\frac{n}{2}-1}(rs) J_{\frac{n}{2}}(ust) ds.$$

Apply (8) to the first Bessel function on the right-hand side. Estimates for the remainder terms are obvious, so we have to estimate

$$\int_0^{\frac{r}{2}} F_\alpha(t) t^{\alpha-\frac{1}{2}} dt \int_0^1 (1-u)^{\alpha-1} u^{\frac{n}{2}-1} du \int_A^\infty \frac{\cos(rs-h)}{\sqrt{s}} J_{\frac{n}{2}}(ust) ds$$

where h is a number. Integration by parts in the integral

$$\int_A^\infty \frac{\cos(rs-h)}{\sqrt{s}} J_{\frac{n}{2}}(ust) ds = \int_A^\infty \frac{\cos(rs-h)}{\sqrt{s} s^{\frac{n}{2}}} s^{\frac{n}{2}} J_{\frac{n}{2}}(ust) ds$$

and simple calculations using (7) and (8) yield that the problem is to show that the integral

$$ut \int_A^\infty J_{\frac{n}{2}-1}(ust) \frac{\sin(rs-h)}{\sqrt{s}} ds$$

tends to zero as $A \rightarrow \infty$. We are going to apply (8) to the Bessel function. For this we had to get the factor ut before the integral. It provides convergence of other integrals, in t and u . Again for the remainder terms in (8) estimates are similar to those afore-mentioned. And for the main term we get the inner integral in the form

$$\sqrt{ut} \int_A^\infty \frac{\cos(ust-l) \sin(rs-h)}{s} ds.$$

Since $ut \leq \frac{r}{2}$ for $t \in [0, \frac{r}{2}]$ and $u \in [0, 1]$, the latter integral is obviously small as $A \rightarrow \infty$. The proof is complete. \square

4. FOURIER TRANSFORM: FROM MANY DIMENSIONS TO ONE DIMENSION.

4.1. It happens sometimes that certain conditions provide that the difference is bounded between an integral of the absolute value of the multiple Fourier transform and an integral of the absolute value of the one-dimensional Fourier transform of some other function. Such a result was obtained by Podkorytov [P], provides that the function under consideration is radial and boundedly supported, and its Fourier transform is integrable. In the following theorem this process is realized under essentially different conditions. This will allow us to generalize some important one-dimensional results to the multi-dimensional case. Besides that, we get an easier way to estimate the growth of the Fourier transform when it is non-integrable.

Theorem 2. *Let a function f be radial and satisfying conditions (1) - (4) with $\alpha = \frac{n-1}{2}$. Suppose further that*

$$(14) \quad \int_0^1 \frac{|F(t)|}{t} dt < \infty.$$

Then for $n \geq 2$

$$\begin{aligned} \int_{1 \leq |x| \leq N} |\hat{f}(x)| dx &= \frac{2^{\frac{n+3}{2}} \pi^n}{\Gamma\left(\frac{n}{2}\right)} \int_1^N \left| \int_0^\infty F(t) \sin\left(rt - \frac{\pi n}{2}\right) dt \right| dr \\ &\quad + \theta \left(V_F + \int_0^1 \frac{|F(t)|}{t} dt \right). \end{aligned}$$

where $|\theta| \leq C$.

Proof. After using (6) and passage to spherical coordinates, multidimensional integration is reduced to the one-dimensional, and one has to estimate

$$\int_1^N r^{\frac{n}{2}} \left| \int_0^\infty F(t) t^{\frac{n}{2}} Q(rt) dt \right| dr.$$

When $rt < 1$, we get, as when proving Lemma 2

$$(rt)^{\frac{n}{2}} Q(rt) = \Gamma\left(\frac{n-1}{2}\right) (rt)^{\frac{1}{2}} J_{n-\frac{3}{2}}(rt) + O(1).$$

Now (9) yields

$$\begin{aligned} \int_{1 \leq |x| \leq N} |\hat{f}(x)| dx &= \int_1^N \left| \frac{2^{\frac{n+3}{2}} \pi^n}{\Gamma\left(\frac{n}{2}\right)} \int_0^{\frac{1}{r}} F(t) \sin\left(rt - \frac{\pi n}{2}\right) dt \right. \\ &\quad \left. + r^{\frac{n}{2}} \int_{\frac{1}{r}}^\infty F(t) t^{\frac{n}{2}} Q(rt) dt \right| dr + \theta \int_1^N \int_0^{\frac{1}{r}} |F(t)| dt dr. \end{aligned}$$

Observe that

$$\int_1^N \int_0^{\frac{1}{r}} |F(t)| dt dr \leq \int_1^\infty \int_0^{\frac{1}{r}} |F(t)| dt dr \leq \int_0^1 \frac{|F(t)|}{t} dt.$$

Take now $rt \geq 1$. Applying Lemma 2 with $p = 2, \mu = \frac{n-3}{2}, \beta = \frac{n}{2} - 1$, we get

$$\begin{aligned} (rt)^{\frac{n}{2}} Q(rt) &= \Gamma\left(\frac{n-1}{2}\right) (rt)^{\frac{1}{2}} J_{n-\frac{3}{2}}(rt) \\ &\quad + \frac{n-3}{2} \Gamma\left(\frac{n+1}{2}\right) (rt)^{-\frac{1}{2}} J_{n-\frac{1}{2}}(rt) + \theta (rt)^{-2}. \end{aligned}$$

The remainder is estimated easily by

$$\begin{aligned} \int_1^N r^{-2} dr \int_{\frac{1}{r}}^\infty |F(t)| t^{-2} dt &\leq \int_0^1 |F(t)| t^{-2} dt \int_{\frac{1}{t}}^\infty r^{-2} dr \\ &\quad + \int_1^\infty |F(t)| t^{-2} dt \int_1^\infty r^{-2} dr \leq \int_1^\infty |F(t)| \frac{dt}{t} + V_F. \end{aligned}$$

Take away the main term of $\Gamma\left(\frac{n-1}{2}\right)(rt)^{\frac{1}{2}}J_{n-\frac{3}{2}}(rt)$ (see (8)) from the asymptotic expansion for $(rt)^{\frac{n}{2}}Q(rt)$ and estimate the integral of the function F multiplied by

$$\begin{aligned} & \Gamma\left(\frac{n-1}{2}\right)(rt)^{\frac{1}{2}}\left[J_{n-\frac{3}{2}}(rt) + \sqrt{\frac{2}{\pi rt}}\sin\left(rt - \frac{\pi n}{2}\right)\right] \\ & + \frac{n-3}{2}\Gamma\left(\frac{n+1}{2}\right)(rt)^{-\frac{1}{2}}J_{n-\frac{1}{2}}(rt). \end{aligned}$$

Integrating by parts, we obtain

$$\begin{aligned} & \int_1^N r^{\frac{1}{2}} \left| F(t) \int_t^\infty s^{\frac{1}{2}} \left[J_{n-\frac{3}{2}}(rs) + \sqrt{\frac{2}{\pi rs}} \sin\left(rs - \frac{\pi n}{2}\right) \right] ds \right|_{\frac{1}{r}}^\infty \\ & - \int_{\frac{1}{r}}^\infty dF(t) \int_t^\infty s^{\frac{1}{2}} \left[J_{n-\frac{3}{2}}(rs) + \sqrt{\frac{2}{\pi rs}} \sin\left(rs - \frac{\pi n}{2}\right) \right] ds \Big| dr \\ & + \int_1^N r^{-\frac{3}{2}} \left| F(t) t^{-\frac{1}{2}} J_{n-\frac{3}{2}}(rt) \right|_{\frac{1}{r}}^\infty + \int_{\frac{1}{r}}^\infty t^{-\frac{1}{2}} J_{n-\frac{3}{2}}(rt) dF(t) \\ & - \frac{1}{2} \int_{\frac{1}{r}}^\infty F(t) t^{-\frac{3}{2}} J_{n-\frac{3}{2}}(rt) dt \Big| dr. \end{aligned}$$

Taking into account $|J_{n-\frac{3}{2}}(rt)| \leq C(rt)^{-\frac{1}{2}}$, the second integral (in r) is estimated by

$$\begin{aligned} & \int_1^N \left| F\left(\frac{1}{r}\right) \right| \frac{dr}{r} + \int_1^N r^{-2} dr \int_{\frac{1}{r}}^\infty \frac{1}{t} |dF(t)| \\ & + \int_1^N r^{-2} dr \int_{\frac{1}{r}}^\infty |F(t)| t^{-2} dt \leq 2 \int_0^1 |F(t)| \frac{dt}{t} \\ & + V_F + \int_0^1 \frac{1}{t} |dF(t)| \int_{\frac{1}{t}}^\infty r^{-2} dr + \int_1^\infty \frac{1}{t} |dF(t)| \int_1^\infty r^{-2} dr \\ & \leq 2 \left(V_F + \int_0^1 |F(t)| \frac{dt}{t} \right). \end{aligned}$$

The integrated term under the sign of absolute value in the first integral in r

obviously vanishes at infinity. For $t = \frac{1}{r}$ use (8) and obtain

$$\begin{aligned} & \left| \int_{\frac{1}{r}}^{\infty} s^{\frac{1}{2}} \left[J_{n-\frac{1}{2}}(rs) - \sqrt{\frac{2}{\pi rs}} \sin \left(rs - \frac{\pi n}{2} \right) \right] ds \right| \\ & \leq Cr^{-\frac{3}{2}} \left| \int_{\frac{1}{r}}^{\infty} \frac{1}{s} \sin \left(rs - \frac{\pi n}{2} \right) ds \right| + Cr^{-\frac{5}{2}} \left| \int_{\frac{1}{r}}^{\infty} s^{-2} ds \right| = O(r^{-\frac{3}{2}}), \end{aligned}$$

and

$$\int_1^N \left| F \left(\frac{1}{r} \right) \right| \frac{dr}{r} \leq \int_0^1 |F(t)| \frac{dt}{t}.$$

Apply again (8) to the integral remained and get

$$\begin{aligned} & \int_1^N \frac{dr}{r} \left| \int_{\frac{1}{r}}^{\infty} dF(t) \int_t^{\infty} \frac{1}{s} \sin \left(rs - \frac{\pi n}{2} \right) ds \right| \\ & + \int_1^N r^{-2} dr \int_{\frac{1}{r}}^{\infty} |dF(t)| \int_t^{\infty} s^{-2} ds \leq 3 \int_1^{\infty} r^{-2} dr \int_{\frac{1}{r}}^{\infty} \frac{1}{t} |dF(t)|. \end{aligned}$$

But the right-hand side was already estimated above. The proof of the theorem is complete. \square

Remark 3. Condition (14) is essential. Let $n = 3$. Consider the function

$$f(x) = \sin \left(\ln \ln \frac{e}{|x|} \right)$$

for $|x| \in [0, 1]$, and 0 otherwise. We have

$$f'_0(t) = \frac{1}{t \ln \frac{e}{t} \ln \ln \frac{e}{t}} \cos \left(\ln \ln \frac{e}{t} \right)$$

and $F(t) = t f'_0(t)$. This function obviously satisfies conditions (1)-(4). It is easy to see that

$$\int_0^1 \frac{|F(t)|}{t} dt = \int_0^1 |f'(t)| dt = \infty.$$

Consider

$$\begin{aligned} \int_{1 \leq |x| \leq N} |\hat{f}(x)| dx &= C \int_1^N r^{\frac{3}{2}} dr \left| \int_0^1 F(t) t^{\frac{3}{2}} dt \int_0^1 s^{\frac{3}{2}} J_{\frac{1}{2}}(rts) ds \right| \\ &= C \sqrt{\frac{2}{\pi}} \int_1^N r dr \left| \int_0^1 F(t) t dt \int_0^1 s \sin rts ds \right| \\ &= C \sqrt{\frac{2}{\pi}} \int_1^N dr \left| \int_0^1 F(t) \cos rt dt - \int_0^1 \frac{1}{t} F(t) \sin rt dt \right|. \end{aligned}$$

It suffices now to prove that

$$\lim_{N \rightarrow \infty} \int_1^N \frac{dr}{r} \left| \int_0^1 \frac{F(t)}{t} \sin rt \, dt \right| = \infty.$$

We have after integrating by parts

$$\int_1^N \frac{dr}{r} \left| \int_0^1 \frac{F(t)}{t} \sin rt \, dt \right| = \int_1^N dr \left| \int_0^1 f_0(t) \cos rt \, dt \right|,$$

and one has to prove that the one-dimensional Fourier transform of the function f_0 is nonintegrable. Indeed, if it were integrable, the following condition would valid necessarily (see e.g., [K], Ch.2, §10): the integral $\int_0^1 \frac{f_0(t)}{t} dt$ converges. But it is easy to see that this integral diverges for our function. Thus, we have built the counterexample which shows that condition (14) is sharp.

Another, somewhat more complicate example suggested by E. Belinskii was given in [L2]:

$$f(x) = \frac{1}{\ln \ln \frac{e}{|x|}} \sin \left(\ln \ln \frac{e}{|x|} \right).$$

It was constructed in order f to be continuous at the origin.

4.2. Let us obtain a generalization, to the multiple case, of the Zygmund-Bochkarev criterion for the absolute convergence of Fourier series of the function of bounded variation (see [B1], [B2], Ch.2, Th.3.1). Let us note that an integral analog of the Stechkin criterion (see e.g. [Ba]) is proved by Trigub ([T4], Th. 2). We use the standard notation ω for the modulus of continuity.

Corollary 1. *Let a radial function be boundedly supported, satisfy conditions (1) and (4), and $F(0) = 0$. Then the condition*

$$(15) \quad \sum_{k=1}^{\infty} \frac{1}{k} \sqrt{\omega \left(F ; \frac{1}{k} \right)} < \infty$$

is sufficient and, on the whole class, necessary for $\hat{f} \in L^1(\mathbb{R}^n)$.

Proof. It is natural to consider those functions which are satisfying $\sum_{k=1}^{\infty} \frac{1}{k} \omega \left(F ; \frac{1}{k} \right) <$

∞ . Since $|F(t)| \leq \omega \left(F ; \frac{1}{k} \right)$, the condition $\sum_{k=1}^{\infty} \frac{1}{k} \left| F\left(\frac{1}{k}\right) \right| < \infty$ provides that (14)

holds. Now we are able to apply Theorem 2 and to reduce the problem to the one-dimensional one (see [B1] or [B2], Ch.2, §3, Th.3.1). True, the absolute convergence of the Fourier series was investigated there, but it is closely connected with the integrability of the Fourier transforms due to the following theorem of Trigub (see [T0], [T1]; an application to summability is given in [T2], Corollary 2):

Theorem A1. *Let $f(t)$ be a boundedly supported function of one variable, and $f_1(t) = tf(t)$. Then $\hat{f} \in L^1(\mathbb{R}^1)$ if and only if the functions f and f_1 after periodic continuation have absolutely convergent Fourier series.*

It remains to note that $\omega \left(F ; \frac{1}{k} \right) \leq C \max\{\omega(F, t) ; t\}$, and the corollary is proved. \square

Example 1. Consider the function $f(x) = (1 - |x|^\alpha)_+^\beta$. It was proved earlier (see [Lf], [T3]) that for $\alpha > 0$ and $\beta > \frac{n-1}{2}$ one has $\hat{f} \in L^1(\mathbb{R}^n)$. Let us establish this fact by means of Corollary 1. Condition (1) is evidently satisfied. The same may be said about conditions (4) and (15) for n odd. For n even to verify (4), it suffices to show that for $\psi(t) = t^\gamma(a - t^\alpha)_+^{\frac{1}{2}+\varepsilon}$, with $\varepsilon \geq 0$, $\gamma \geq 0$, the function $t^{\frac{1}{2}}\psi^{(\frac{1}{2})}(t)$ is of bounded variation. We have

$$\begin{aligned} \psi^{(\frac{1}{2})}(t) &= \frac{d}{dt} \int_t^1 (s-t)^{-\frac{1}{2}} s^\gamma (1-s^\alpha)^{\frac{1}{2}+\varepsilon} ds \\ &= \int_t^1 (s-t)^{-\frac{1}{2}} \frac{d}{ds} \left\{ s^\gamma (1-s^\alpha)^{\frac{1}{2}+\varepsilon} \right\} ds. \end{aligned}$$

This means that the boundedness of variation of the function

$$t^{\frac{1}{2}} \int_t^1 (s-t)^{-\frac{1}{2}} s^{\zeta-1} (1-s^\alpha)^{\varepsilon-\frac{1}{2}} ds,$$

with $\varepsilon \geq 0$, $\zeta > 0$, should be established. Further,

$$s^{\zeta-1} (1-s^\alpha)^{\varepsilon-\frac{1}{2}} = (1-s^\alpha)^{\varepsilon-\frac{1}{2}} (s^{\zeta-1} - 1) + (1-s^\alpha)^{\varepsilon-\frac{1}{2}}.$$

Denoting $C_\alpha = \lim_{s \rightarrow 1} \left(\frac{1-s^\alpha}{1-s} \right)^{\varepsilon-\frac{1}{2}}$, we have

$$\begin{aligned} (1-s^\alpha)^{\varepsilon-\frac{1}{2}} &= (1-s^\alpha)^{\varepsilon-\frac{1}{2}} - C_\alpha (1-s)^{\varepsilon-\frac{1}{2}} + C_\alpha (1-s)^{\varepsilon-\frac{1}{2}} \\ &= (1-s^\alpha)^{\varepsilon+\frac{1}{2}} \frac{\left(\frac{1-s^\alpha}{1-s} \right)^{\varepsilon-\frac{1}{2}} - C_\alpha}{1-s} + C_\alpha (1-s)^{\varepsilon-\frac{1}{2}}. \end{aligned}$$

Thus, the boundedness of variation of the function

$$t^{\frac{1}{2}} \int_t^1 (s-t)^{-\frac{1}{2}} (1-s)^{\varepsilon-\frac{1}{2}} ds$$

should be established. But, after a simple change of variables, it is equal to the function

$$t^{\frac{1}{2}} (1-t)^\varepsilon \int_0^1 s^{-\frac{1}{2}} (1-s)^{\varepsilon-\frac{1}{2}} ds,$$

and (4) is now obvious. Moreover, this makes (14) obvious too, and this completes the proof of the example.

Example 2. Consider $f(x) = \frac{1 - (1 - |x|^\alpha)_+^\beta}{|x|^r}$ with $\alpha > r$ and $\beta > \frac{n-1}{2}$. Let us show again that $\hat{f} \in L^1(\mathbb{R}^n)$. On $[0, 1]$ the argument from Example 1 is applicable. Hence for $t \in [0, 1]$ we have

$$F(t) = C_1 t^{\alpha-r} (1-t)^{\beta-\frac{n-1}{2}} + g(t),$$

where g is a continuously differentiable function, $g(0) = 0$. It is clear that (14) is satisfied. Applying, if needed, the formula of fractional derivation (see [BE2], p. 201), we obtain $t^{\frac{n-1}{2}} (t^{-r})^{(\frac{n-1}{2})} = C_2 t^{-r}$. Use now Theorem 2 and integrate by parts in the one-dimensional integral. After that we have to estimate the following value:

$$\int_1^N \frac{1}{s} \left| \int_0^\infty F'(t) \cos\left(st - \frac{\pi n}{2}\right) dt \right| ds.$$

On $[0, 1]$ we have $F'(t) \in \text{Lip } \varepsilon$ in L^1 metric, for some $\varepsilon > 0$. But the integral $\int_1^N \frac{ds}{s^{1+\varepsilon}}$ converges. For $t \in [1, \infty]$, we obtain by integration by parts:

$$\begin{aligned} \int_1^N \frac{1}{s} \left| \int_1^\infty t^{-1-r} \cos\left(st - \frac{\pi n}{2}\right) dt \right| ds &= \int_1^N \frac{1}{s} \left| \frac{1}{s} t^{-1-r} \sin\left(st - \frac{\pi n}{2}\right) \right|_1^\infty \\ &+ \frac{1+r}{s} \int_1^\infty t^{-2-r} \left| \sin\left(st - \frac{\pi n}{2}\right) \right| dt \Big| ds \leq C \int_1^N \frac{ds}{s^2}, \end{aligned}$$

and Example 2 is proved.

5. RADIAL FUNCTIONS WITH CONVEXITY CONDITIONS.

5.1. An asymptotics of the Fourier transform may be established not very often. Some special conditions, like convexity, must be layed on a function. G. E. Shilov (see e.g., [Ba], p. 632) was the first who studied the Fourier coefficients of convex functions. In [T2] an asymptotics of the Fourier transform of a convex function is given in a sharp form as follows:

Theorem A2. *If f is convex on $[a, b]$, where $-\infty < a < b \leq +\infty$, and $|f'(b)| < \infty$, then for each $r \in \mathbb{R}^1$, $|r| \geq 2$,*

$$\int_a^b f(t) e^{-irt} dt = \frac{i}{r} \left\{ f(b) e^{-ibr} - f\left(a + \frac{d}{|r|}\right) e^{-iav} \right\} + \theta \gamma(|r|),$$

where $d = \min\{b - a, \pi\}$, $|\theta| \leq C$, and γ is monotone decreasing so that

$$\int_2^\infty \gamma(t) dt \leq \frac{1}{d} V_f + |f'(b)|.$$

5.2. In [T3], [T4], [L1], this theorem was generalized to the radial case. Here we prove the stronger result formulated in [BL1]

Theorem 3. *Let a radial function f satisfy (1). Suppose further that f_0 is supported on $[0, 1]$ and F is continuous on $[0, \infty)$ and convex on $[0, 1]$. Then for $|x| = r \geq 2$*

$$\hat{f}(x) = 2^{\frac{n+1}{3}} \pi^{\frac{n-1}{2}} (-1)^{[\frac{n}{2}]} r^{-n} F\left(1 - \frac{1}{r}\right) \cos\left(r - \frac{\pi n}{2}\right) + \theta\gamma(r),$$

where $|\theta| \leq C$ and γ decreases monotonously so that

$$\int_2^\infty r^{n-1} |\gamma(r)| dr \leq V_F.$$

Proof. Using (6) we obtain

$$\begin{aligned} \hat{f}(x) &= \hat{f}_0(r) = \frac{(2\pi)^{\frac{n}{2}} (-1)^{[\frac{n}{2}]} r^{1-\frac{n}{2}}}{\Gamma\left(\frac{n-1}{2}\right)} \int_0^1 F(t) t^{\frac{n}{2}} Q(rt) dt \\ &= \frac{(2\pi)^{\frac{n}{2}} (-1)^{[\frac{n}{2}]} r^{1-\frac{n}{2}}}{\Gamma\left(\frac{n-1}{2}\right)} \left\{ F\left(1 - \frac{1}{r}\right) \int_0^1 t^{\frac{n}{2}} Q(rt) dt \right. \\ &\quad + \int_{1-\frac{1}{r}}^1 \left[F(t) - F\left(1 - \frac{1}{r}\right) \right] t^{\frac{n}{2}} Q(rt) dt \\ &\quad \left. - \frac{1}{r} \int_{\frac{1}{r}}^{1-\frac{1}{r}} F'(t) t^{\frac{n}{2}} q(rt) dt + \int_0^{\frac{1}{r}} \left[F(t) - F\left(\frac{1}{r}\right) \right] t^{\frac{n}{2}} Q(rt) dt \right\}. \end{aligned}$$

Lemma 1 and (7) yield

$$\begin{aligned} (16) \quad & \int_0^1 t^{\frac{n}{2}} Q(rt) dt = \frac{1}{r} q(r) \\ &= \Gamma\left(\frac{n-1}{2}\right) r^{-\frac{n+1}{2}} J_{n-\frac{1}{2}}(r) + \zeta_n r^{-\frac{n+2}{2}} + O\left(r^{-\frac{n+4}{2}}\right). \end{aligned}$$

We denoted $\zeta_{\frac{n-1}{2}, n}$ by ζ_n . The remainder term may be obviously referred to as γ . In order to obtain the main term of asymptotics apply (8) to $J_{n-\frac{1}{2}}(r)$. Then, using Lemma 2 for the upper estimate, we obtain

$$\begin{aligned} (17) \quad & \left| \int_{1-\frac{1}{r}}^1 [F(t) - F(1 - \frac{1}{r})] t^{\frac{n}{2}} Q(rt) dt \right| = \left| \int_{1-\frac{1}{r}}^1 t^{\frac{n}{2}} Q(rt) dt \int_{1-\frac{1}{r}}^t F'(s) ds \right| \\ &= \left| \int_{1-\frac{1}{r}}^1 F'(s) ds \int_{1-\frac{1}{r}}^1 t^{\frac{n}{2}} Q(rt) dt \right| \leq C r^{-\frac{n}{2}} \int_{1-\frac{1}{r}}^1 |F'(s)| (1-s) ds. \end{aligned}$$

Let us verify that the right-hand side satisfies conditions which are defining γ . Monotonicity is obvious. Further

$$\int_2^\infty r^{n-1} r^{1-\frac{n}{2}} r^{-\frac{n}{2}} dr \int_{1-\frac{2}{r}}^1 |F'(s)|(1-s) ds = \int_0^1 |F'(s)|(1-s) ds \int_2^{\frac{1}{1-s}} dr \leq 2V_F.$$

Lemma 1 yields for large rt

$$t^{\frac{n}{2}} q(rt) = \Gamma\left(\frac{n-1}{2}\right) t^{\frac{1}{2}} r^{-\frac{n-1}{2}} J_{n-\frac{1}{2}}(rt) + \zeta_n r^{-\frac{n}{2}} + O\left(\frac{1}{t} r^{-\frac{n+2}{2}}\right).$$

For rt small the remainder term is $O(t^{\frac{n}{2}})$. Then

$$\begin{aligned} & -\frac{1}{r} \int_{\frac{1}{r}}^{1-\frac{1}{r}} F'(t) t^{\frac{n}{2}} q(rt) dt + \int_0^{\frac{1}{r}} \left[F(t) - F\left(\frac{1}{r}\right) \right] t^{\frac{n}{2}} Q(rt) dt \\ & = -\Gamma\left(\frac{n-1}{2}\right) r^{-\frac{n+1}{2}} \int_{\frac{1}{r}}^{1-\frac{1}{r}} F'(t) t^{\frac{1}{2}} J_{n-\frac{1}{2}}(rt) dt - \zeta_n r^{-\frac{n+2}{2}} F\left(1 - \frac{1}{r}\right) \\ (18) \quad & + \zeta_n r^{-\frac{n+2}{2}} F\left(\frac{1}{r}\right) + O\left(r^{-\frac{n+4}{2}} \int_{\frac{1}{r}}^{1-\frac{1}{r}} |F'(t)| \frac{dt}{t}\right) - \frac{1}{r} \int_0^{\frac{1}{r}} F'(t) t^{\frac{n}{2}} q(rt) dt. \end{aligned}$$

But we have

$$\begin{aligned} & -\frac{1}{r} \int_0^{\frac{1}{r}} F'(t) t^{\frac{n}{2}} q(rt) dt = -\Gamma\left(\frac{n-1}{2}\right) r^{-\frac{n+1}{2}} \int_0^{\frac{1}{r}} F'(t) t^{\frac{1}{2}} J_{n-\frac{1}{2}}(rt) dt \\ & - \zeta_n r^{-\frac{n+2}{2}} F\left(\frac{1}{r}\right) + O\left(\frac{1}{r} \int_0^{\frac{1}{r}} |F'(t)| t^{\frac{n}{2}} dt\right). \end{aligned}$$

Now the second summand annihilates the third summand on the right-hand side of (18), and after applying (8) the first summand and the remainder term are estimated by

$$(19) \quad r^{-\frac{n}{2}} \int_0^{\frac{2}{r}} |F'(t)| t dt.$$

Integrability of (19) over $[2, \infty]$ with the weight $r^{1-\frac{n}{2}} r^{n-1}$ is verified like for (17). The second summands in (16) and (18) are mutually annihilated as well. Let us prove that the remainder term in (18) is a part of γ .

$$\begin{aligned} & \int_2^\infty r^{n-1} dr r^{-n-1} \int_{\frac{1}{r}}^{1-\frac{1}{r}} |F'(t)| \frac{dt}{t} = \int_2^\infty \frac{dr}{r^2} \left(\int_{\frac{1}{r}}^{\frac{1}{2}} + \int_{\frac{1}{2}}^{1-\frac{1}{r}} \right) |F'(t)| \frac{dt}{t} \\ & \leq \int_{\frac{1}{2}}^{\frac{1}{2}} |F'(t)| \frac{dt}{t} \int_2^\infty \frac{dr}{r^2} + 2 \int_{\frac{1}{2}}^1 |F'(t)| dt \leq 2 \int_{\frac{1}{2}}^1 |F'(t)| dt. \end{aligned}$$

It remains to estimate the first summand on the right-hand side of (18). Integrating by parts we obtain

$$\begin{aligned} \int_{\frac{1}{r}}^{1-\frac{1}{r}} F'(t) t^{\frac{1}{2}} J_{n-\frac{1}{2}}(rt) dt &= F'(t) \int_0^t s^{\frac{1}{2}} J_{n-\frac{1}{2}}(rs) ds \Big|_{\frac{1}{r}}^{1-\frac{1}{r}} \\ &\quad - \int_{\frac{1}{r}}^{1-\frac{1}{r}} \left[\int_0^t s^{\frac{1}{2}} J_{n-\frac{1}{2}}(rs) ds \right] dF'(t). \end{aligned}$$

Let us estimate the integrated terms. The integral is estimated by Lemma 3. Since F' can be assumed monotone the integrated term for $t = 1 - \frac{1}{r}$ is estimated by (17), and for $t = \frac{1}{r}$ by (19). Again, using Lemma 3 we have

$$\left| \int_{\frac{1}{r}}^{1-\frac{1}{r}} \int_0^t s^{\frac{1}{2}} J_{n-\frac{1}{2}}(rs) ds dF'(t) \right| \leq C \frac{1}{r} \int_{\frac{1}{r}}^{\frac{1}{2}} t^{\frac{1}{2}} |dF'(t)| + Cr^{-\frac{1}{2}} \int_{\frac{1}{2}}^{1-\frac{1}{r}} |dF'(t)|.$$

Since $F'(t)$ is monotone, then

$$\int_{\frac{1}{r}}^{\frac{1}{2}} t^{\frac{1}{2}} |dF'(t)| = \left| \int_{\frac{1}{r}}^{\frac{1}{2}} t^{\frac{1}{2}} dF'(t) \right| = \left| F'(t) t^{\frac{1}{2}} \Big|_{\frac{1}{r}}^{\frac{1}{2}} - \frac{1}{2} \int_{\frac{1}{r}}^{\frac{1}{2}} t^{-\frac{1}{2}} F'(t) dt \right|$$

and

$$\int_{\frac{1}{2}}^{1-\frac{1}{r}} |dF'(t)| = \left| F'(1 - \frac{1}{r}) - F' \left(\frac{1}{2} \right) \right|.$$

Again the integrated terms are estimated by (17), (19), and also by

$$\left| F' \left(\frac{1}{2} \right) \right| \int_2^\infty r^{n-1} r^{-n-\frac{1}{2}} dr = \sqrt{2} \left| F' \left(\frac{1}{2} \right) \right|.$$

Taking into account the factor $r^{-n-\frac{1}{2}}$ we have for the integral the final estimate

$$r^{-n-\frac{1}{2}} \int_{\frac{1}{r}}^{\frac{1}{2}} t^{-\frac{1}{2}} |F'(t)| dt.$$

It remains to verify that this value is also suitable for γ . Again the monotonicity is obvious. Further

$$\begin{aligned} &\int_2^\infty r^{n-1} r^{-n-\frac{1}{2}} dr \int_{\frac{1}{r}}^{\frac{1}{2}} t^{-\frac{1}{2}} |F'(t)| dt \\ &= \int_0^{\frac{1}{2}} |F'(t)| t^{-\frac{1}{2}} dt \int_{\frac{1}{2}}^\infty r^{-\frac{3}{2}} dr = 2 \int_0^{\frac{1}{2}} |F'(t)| dt \leq 2V_F. \end{aligned}$$

The theorem is completely proved. \square

Remark 4. This theorem also shows that convexity conditions posed on F allow to avoid condition (14).

5.3. A criterion of integrability of Fourier transforms is a simple corollary to this theorem.

Corollary 2. *Under conditions of Theorem 3 the Fourier transform \hat{f} is integrable if and only if the integral*

$$(20) \quad \int_0^1 \frac{f_0(t)}{(1-t)^{\frac{n+1}{2}}} dt$$

converges.

Proof. It is obvious that $\hat{f} \in L^1(\mathbb{R}^n)$ is equivalent to the finiteness of

$$\int_0^\infty r^{n-1} r^{-n} \left| F\left(1 - \frac{1}{r}\right) \right| dr = \int_0^{\frac{1}{2}} |F(1-t)| \frac{dt}{t}.$$

Recalling the definition of F and taking into account that by convexity $F(1-t)$ preserves a sign near the origin, we get the following integral to estimate:

$$\int_0^{\frac{1}{2}} \frac{1}{t} f_0\left(\frac{n-1}{2}\right) (1-t) dt.$$

If n is odd the usual integration by parts leads to the required result. When n is even it suffices to consider

$$\int_0^{\frac{1}{2}} t^{-\frac{n}{2}} f_0\left(\frac{1}{2}\right) (1-t) dt = \int_0^\infty t^{-\frac{n}{2}} \frac{d}{dt} \int_{1-t}^1 (s-1+t)^{-\frac{1}{2}} f_0(s) ds dt.$$

Integrate by parts in t . We get

$$\begin{aligned} \int_0^{\frac{1}{2}} t^{-\frac{n}{2}} f_0\left(\frac{1}{2}\right) (1-t) dt &= t^{-\frac{n}{2}} \int_{1-t}^1 (s-1+t)^{-\frac{1}{2}} f_0(s) ds \Big|_0^{\frac{1}{2}} \\ &\quad + \frac{n}{2} \int_0^{\frac{1}{2}} t^{-1-\frac{n}{2}} \int_{1-t}^1 (s-1+t)^{-\frac{1}{2}} f_0(s) ds dt. \end{aligned}$$

The integrated terms are bounded because the integral (20) converges. After changing the order of integration, the last integral is equal to

$$\begin{aligned} &\int_{\frac{1}{2}}^1 f_0(s) ds \int_{1-s}^{\frac{1}{2}} (s-1+t)^{-\frac{1}{2}} t^{-1-\frac{n}{2}} dt \\ &= \int_{\frac{1}{2}}^1 f_0(s) ds \int_{1-s}^\infty (s-1+t)^{-\frac{1}{2}} t^{-1-\frac{n}{2}} dt - \int_{\frac{1}{2}}^1 f_0(s) ds \int_{\frac{1}{2}}^\infty (s-1+t)^{-\frac{1}{2}} t^{-1-\frac{n}{2}} dt. \end{aligned}$$

The last integral is obviously bounded. As in the first one the inner integral is the Weyl integral of order $\frac{1}{2}$ of the function $t^{-1-\frac{n}{2}}$. By the formula in [BE2], p.201 it is equal to $\frac{\Gamma(\frac{n+1}{2})}{\Gamma(\frac{n+2}{2})} (1-t)^{-\frac{n+1}{2}}$, and the corollary is established. \square

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